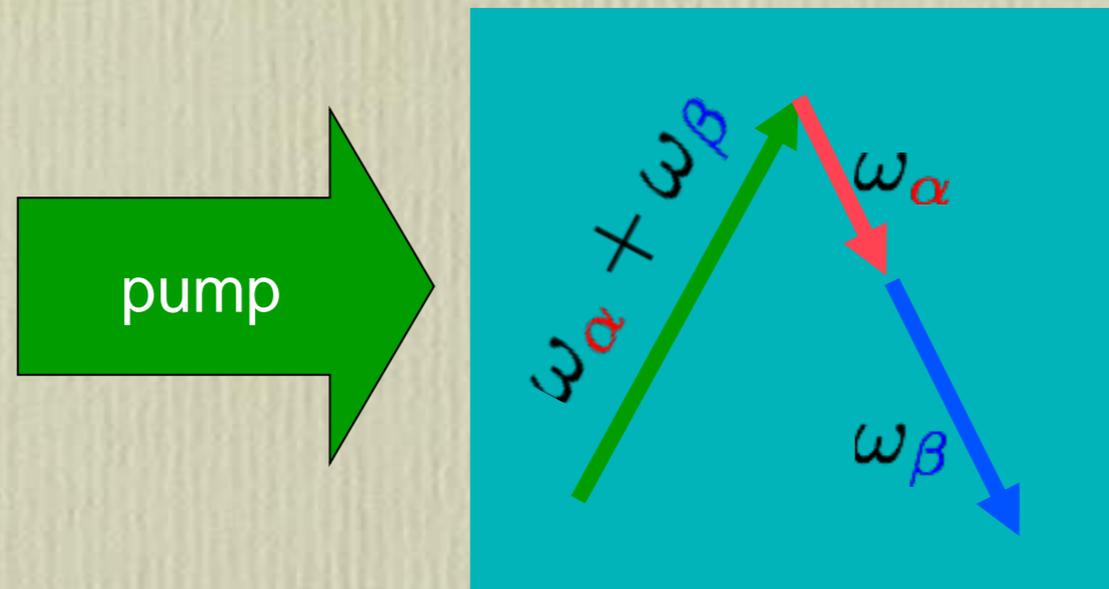


The correlations present in a separable state
turn out to be as efficient
as the correlations arising from entanglement
to boost the two-atom two-photon excitation

two-photon state
produced by parametric down-conversion:



same conclusion:

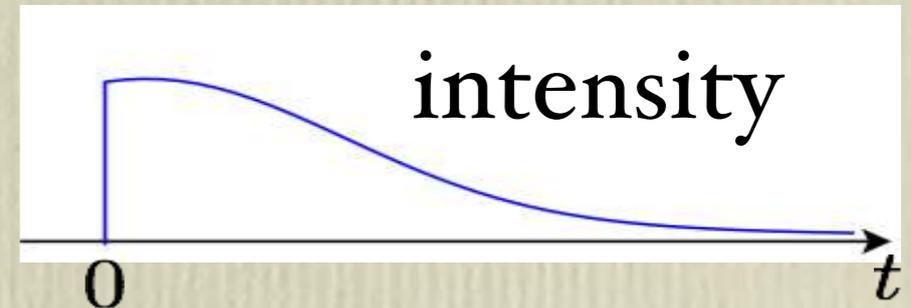
entangled pure state

and correlated separable mixed state

lead to the same enhancement and same resonant character
of the two-atom excitation

two-photon state pulse starting at a given time

$$t=0 \quad |\Psi\rangle = \sum_{kq} c_{kq} |1 : \omega_k; 1 : \omega_q\rangle$$



with mean-field non-zero only for $t > 0$

transition probability is the same, within a factor 2,
for the pure entangled and for the mixed separable state

enhancement factor R with respect to uncorrelated photons:

$$R = \frac{|c(\omega_1, \omega_2)|^2 + |c(\omega_2, \omega_1)|^2}{\sum_{mn} [|c(\omega_1, \omega_n)c(\omega_m, \omega_2)|^2 + |c(\omega_2, \omega_n)c(\omega_m, \omega_1)|^2]}$$

superpositions of coherent states

$$|QMC\rangle = \sum_{kq} c(\omega_k, \omega_q) |\alpha : \omega_k; \alpha : \omega_q\rangle$$

QMC pure state is highly non-classical !

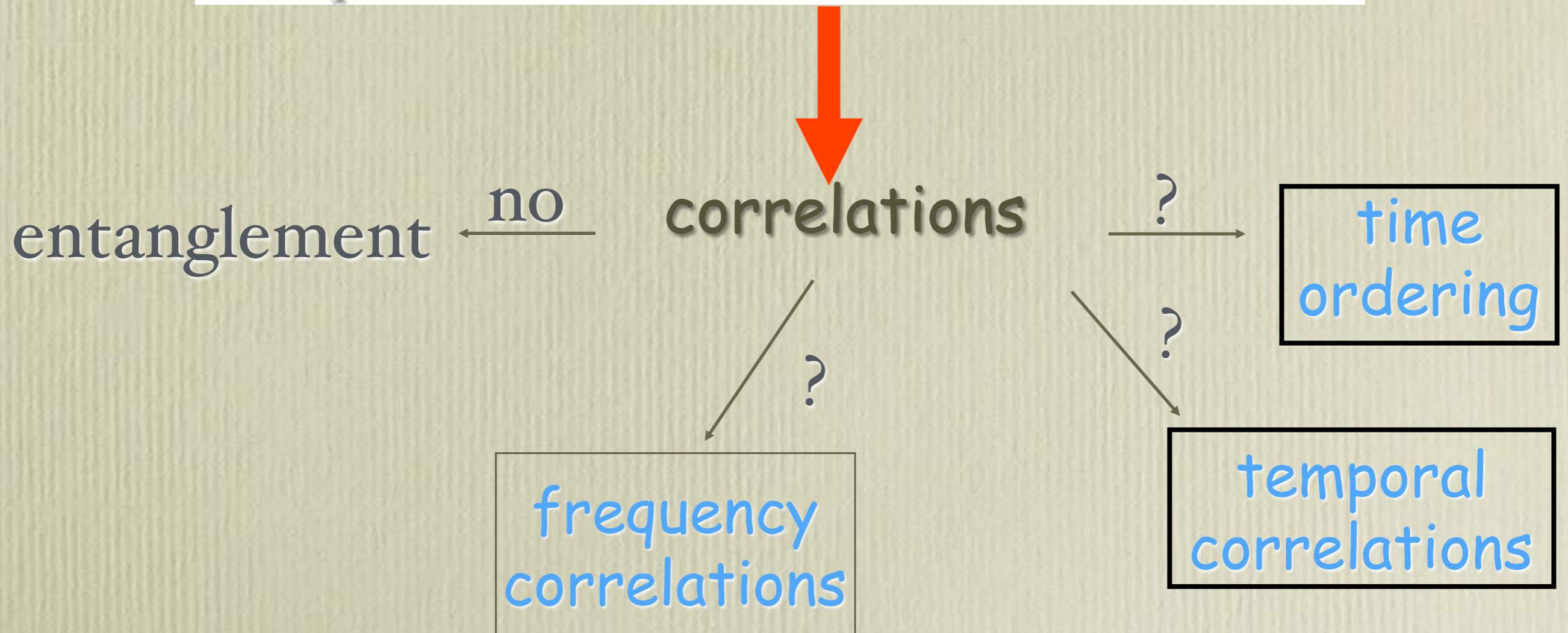
$$\rho_{CMC} = \sum_{kq} |c(\omega_k, \omega_q)|^2 |\alpha : \omega_k; \alpha : \omega_q\rangle \langle \alpha : \omega_k; \alpha : \omega_q|$$

QMC separable mixed state
can be produced by classical means !

$$P_{CMC} \simeq P_{QMC} \simeq |\alpha|^4 P(t)$$

transition probability greatly enhanced

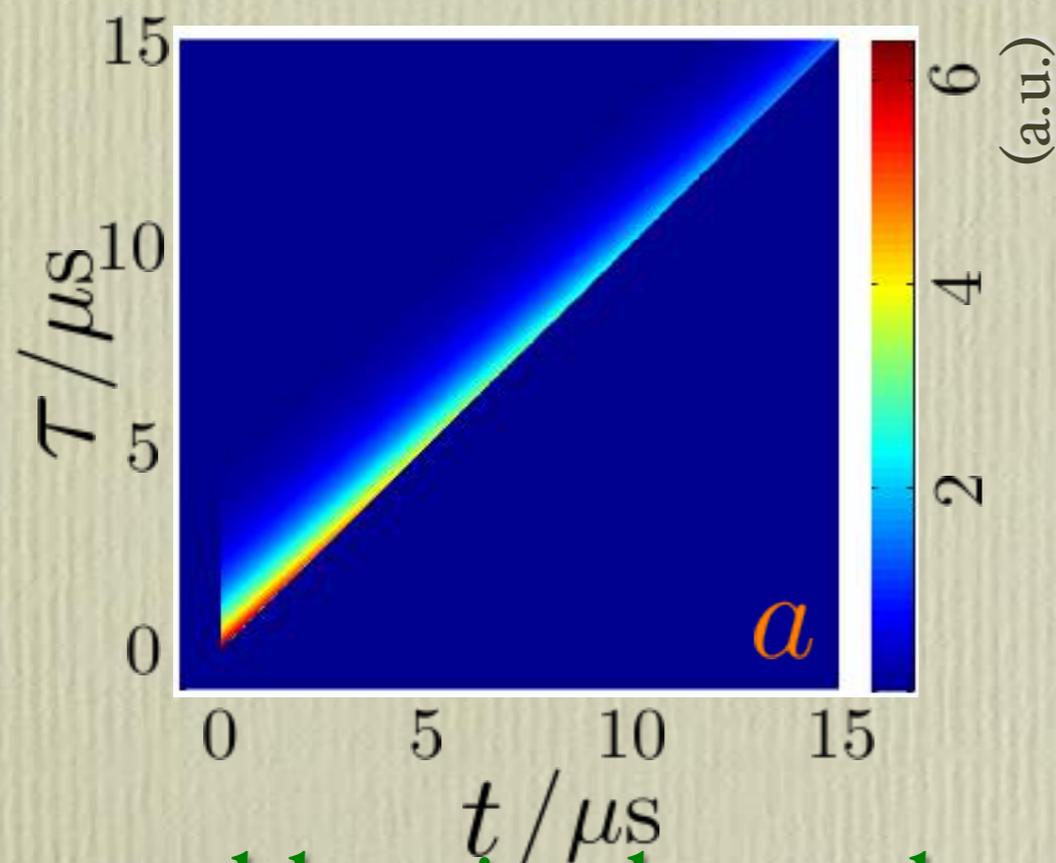
What is the physical origin of the two-photon two-atom enhancement effect?



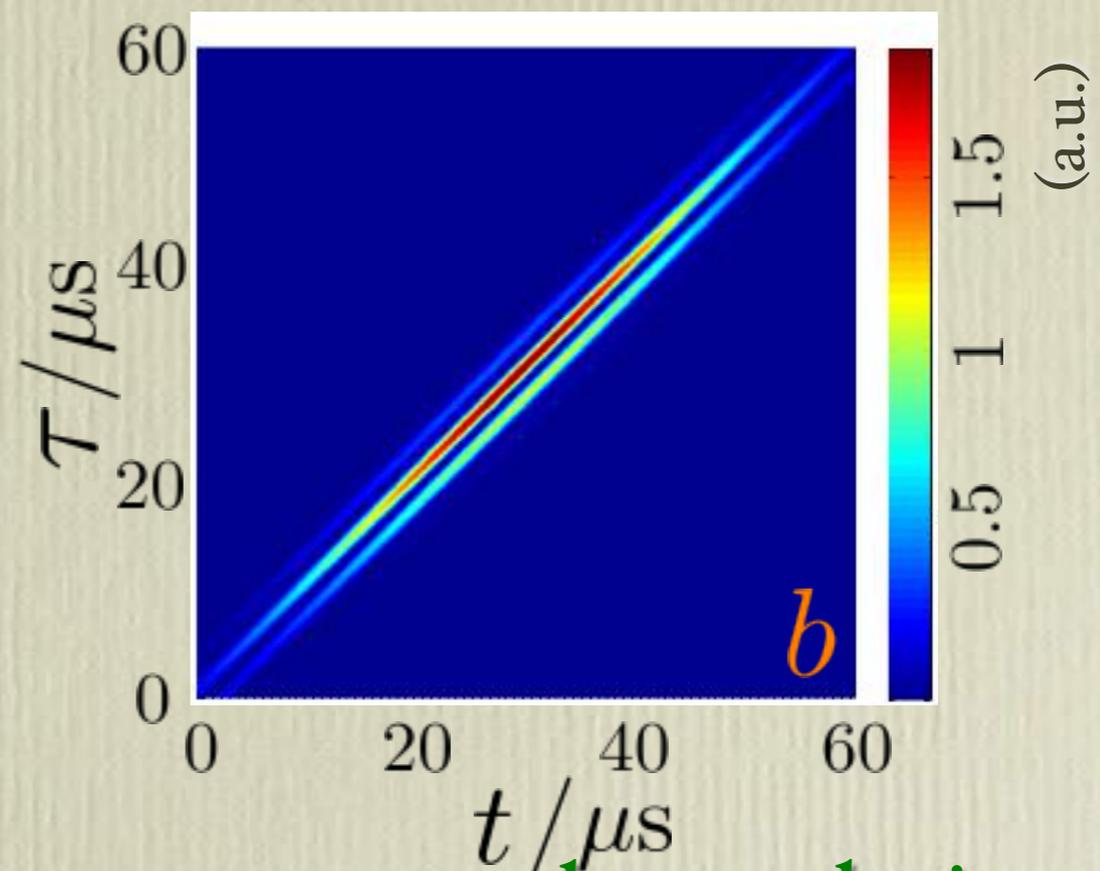
time ordering of the two-photons ?

two-time correlation function

pure cascade state
(time ordered photons)



pure SPDC state
(no time ordering)

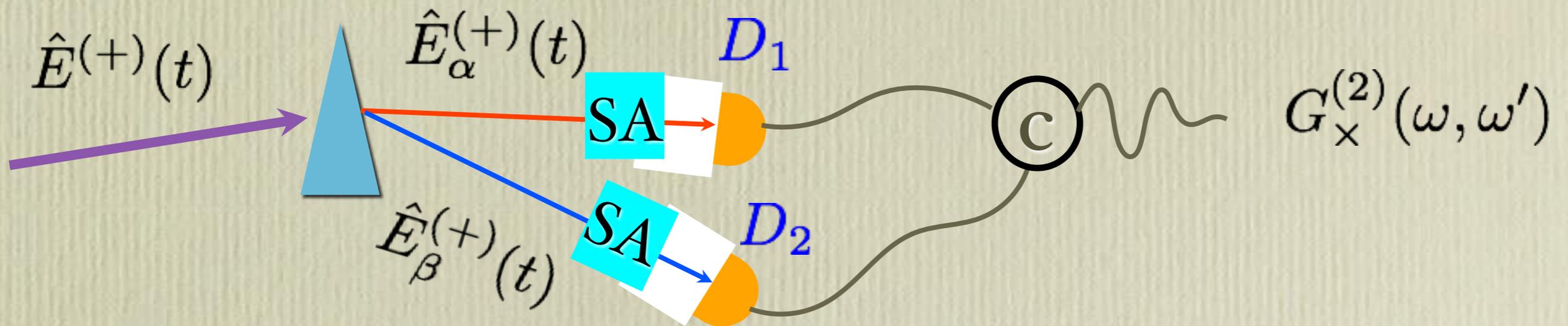


separable mixed state have no temporal correlation !

enhanced proba

time ordering is not at the origin of the effect

frequency correlations between the two-photons ?



SA: spectral analysis

$$g_x^{(2)}(\omega, \omega') = \frac{\langle \hat{\mathcal{E}}_\alpha^{(-)}(\omega) \hat{\mathcal{E}}_\beta^{(-)}(\omega') \hat{\mathcal{E}}_\beta^{(+)}(\omega') \hat{\mathcal{E}}_\alpha^{(+)}(\omega) \rangle}{\langle \hat{\mathcal{E}}_\alpha^{(-)}(\omega) \hat{\mathcal{E}}_\alpha^{(+)}(\omega) \rangle \langle \hat{\mathcal{E}}_\beta^{(-)}(\omega') \hat{\mathcal{E}}_\beta^{(+)}(\omega') \rangle}$$

pure state and mixed state

have the same frequency correlation function

$G_X^{(2)}(\omega, \omega')$ functions

time-frequency entanglement

pure

correlated-separable

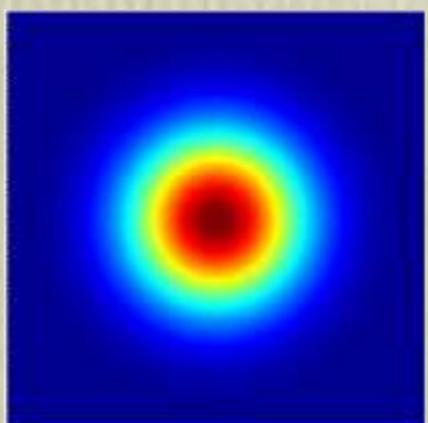
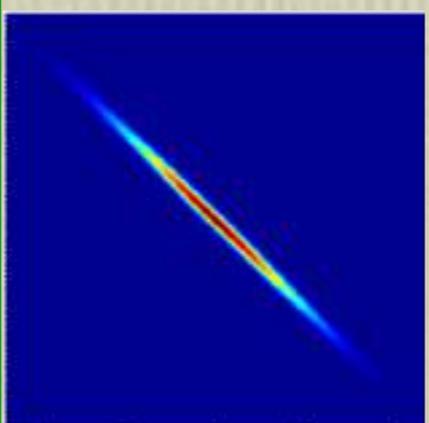
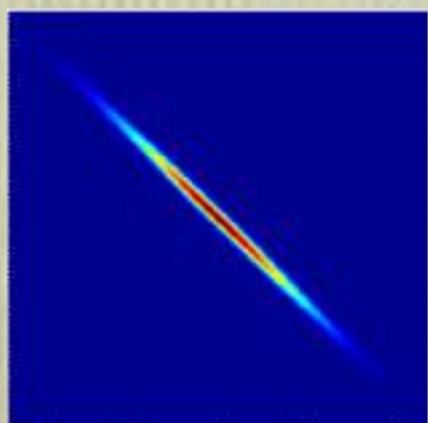
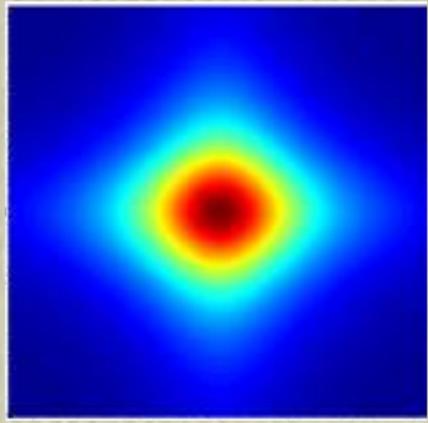
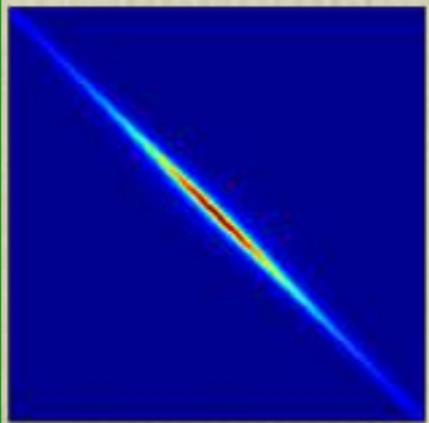
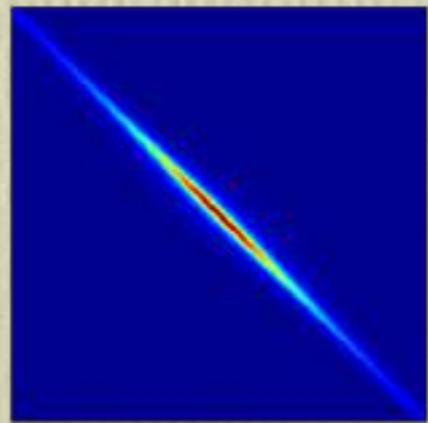
dephased-factorable

$$\Gamma = \gamma_\beta, \sigma = \gamma_\alpha$$

100 200 300

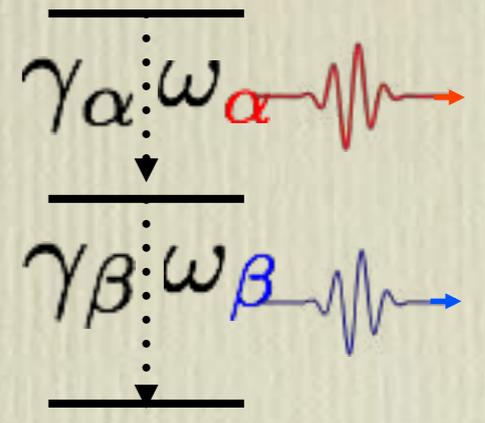
100 200 300

5 10 15 20 25

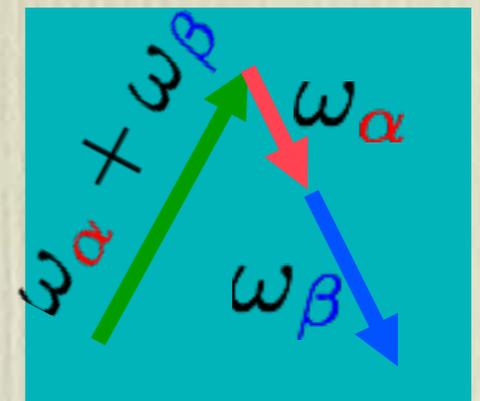


1
0
-1
1
0
-1

cascade



pdc



enhanced proba

$(\omega' - \omega_\beta)/2\gamma_\beta$

-1 0 1 -1 0 1 -1 0 1

The enhancement effect is due to strong frequency anti-correlations in the two-photon state, not time/frequency entanglement

origin of enhancement



yes!

correlations

no

no

entanglement

time ordering



yes!

frequency
(anti-correlations)



no

temporal correlations



?

classical ?



?

quantum ?

Are the required frequency correlations of quantum origin ?

- no quantum discord in the state necessary
- when the cooperativity (Schmidt number)
is high, the enhancement effect is big
- most light states used are not classical

except the separable mix of coherent states

$$\rho_{CMC} = \sum_{kq} |c(\omega_k, \omega_q)|^2 |\alpha : \omega_k; \alpha : \omega_q\rangle \langle \alpha : \omega_k; \alpha : \omega_q|$$

Conclusion

The frequency anti-correlations required for the two-photon two atom enhancement effect

-may be due to the presence of entanglement

-may be due to correlations not related to entanglement,

which can be either of quantum or classical origin